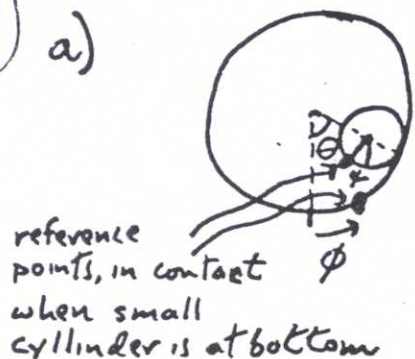


Physics 106a
Final Exam Solutions
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①

① a)



Rolls without slipping $\Rightarrow$

$$r(\theta + \phi) = R(\theta - \phi); \quad \dot{\phi} = \alpha t$$

$$\Rightarrow r\dot{\phi} = (R-r)\dot{\theta} - R\alpha t$$

$$K.E. = \frac{1}{2}m(R-r)^2\dot{\theta}^2 + \frac{1}{4}mr^2\dot{\phi}^2$$

$$= \frac{3}{4}m(R-r)^2\dot{\theta}^2 - \frac{1}{2}m\alpha t R(R-r)\dot{\theta} + \frac{1}{4}mR^2\alpha^2 t^2$$

b) $L = T + mg(R-r)\cos\theta$

$$\Rightarrow \frac{\partial L}{\partial \theta} = -mg(R-r)\sin\theta = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} = \frac{d}{dt} \left[\frac{3}{2}m(R-r)^2\dot{\theta} - \frac{1}{2}mR(R-r)\alpha t \right]$$

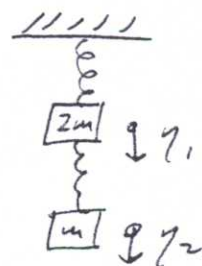
$$\Rightarrow \frac{3}{2}(R-r)\dot{\theta} = \frac{1}{2}R\alpha - g\sin\theta$$

c) $\ddot{\theta} = 0 \Rightarrow \sin\theta = R\alpha/2g; \quad \alpha < 2g/R$

② a) $T = \frac{1}{2}(2m)\dot{\eta}_1^2 + \frac{1}{2}m\dot{\eta}_2^2$

Since gravitational potential energy depends linearly on height of the blocks, gravity does not contribute to quadratic part of V .

$$V = \frac{1}{2}K[\eta_1^2 + (\eta_2 - \eta_1)^2], \quad L = T - V$$



b) we have $\underline{T} = \begin{pmatrix} 2m & 0 \\ 0 & m \end{pmatrix}$ and $\underline{V} = \begin{pmatrix} 2K & -K \\ -K & K \end{pmatrix}$.

so

$$\det(\omega^2 \underline{T} - \underline{V}) = \det \begin{pmatrix} 2m\omega^2 - 2K & K \\ K & m\omega^2 - K \end{pmatrix}$$

$$= 2(m\omega^2 - K)^2 - K^2 = 2m^2\omega^4 - 4Km\omega^2 + K^2 = 0.$$

Thus, $\omega^2 = \frac{1}{4m^2} [4Km \pm \sqrt{16K^2m^2 - 8K^2m^2}] = \frac{K}{m} (1 \pm \frac{1}{\sqrt{2}})$

(2)

c) Let $\omega_1^2 = k/m(1 + \frac{1}{\sqrt{2}})$. Then $\omega_1^2 \underline{I} - \underline{V} = k \begin{pmatrix} \sqrt{2} & 1 \\ 1 & 1/\sqrt{2} \end{pmatrix}$,
 and $(\omega_1^2 \underline{I} - \underline{V}) \vec{X}^{(1)} = 0 \Rightarrow \vec{X}^{(1)} \propto \begin{pmatrix} 1 \\ -\sqrt{2} \end{pmatrix}$.

Normalizing: $\vec{X}^{(1)} = \frac{1}{\sqrt{4m}} \begin{pmatrix} 1 \\ -\sqrt{2} \end{pmatrix}$

Similarly $\omega_2^2 = k/m(1 - \frac{1}{\sqrt{2}}) \Rightarrow \omega_2^2 \underline{I} - \underline{V} = k \begin{pmatrix} -\sqrt{2} & 1 \\ 1 & -1/\sqrt{2} \end{pmatrix}$,
 so $(\omega_2^2 \underline{I} - \underline{V}) \vec{X}^{(2)} = 0 \Rightarrow \vec{X}^{(2)} \propto \begin{pmatrix} 1 \\ \sqrt{2} \end{pmatrix}$.

Normalizing: $\vec{X}^{(2)} = \frac{1}{\sqrt{4m}} \begin{pmatrix} 1 \\ \sqrt{2} \end{pmatrix}$

③ a) $[Q, Q] = [P, P] = 0$; $[Q, P] = \cos^2 \alpha [q, p] - \sin^2 \alpha [p, q] = [q, p] = 1$

b) $P = \frac{1}{\cos \alpha} [p - \sin \alpha q] = \partial G_2 / \partial q$

$Q = \frac{1}{\cos \alpha} [q - \sin \alpha p] = \partial G_2 / \partial P$

$\Rightarrow G_2(q, P) = \frac{1}{\cos \alpha} [qP - \frac{1}{2} \sin \alpha (P^2 + q^2)]$

c) $G_2 \sim qP - \frac{1}{2} \alpha (P^2 + q^2) + O(\alpha^2)$
 $\Rightarrow p^2 + q^2$ conserved

④ a) T is a generating function of the third type $\Rightarrow$

$q = -\frac{\partial T}{\partial p}$ and $\bar{H}(Q, P, t) = H(q, p, t) + \frac{\partial T}{\partial t}$

We require $\bar{H} = 0$ or

$\frac{\partial T}{\partial t} + H(-\frac{\partial T}{\partial p}, p, t) = 0$

b) $H = \frac{1}{2m} p^2 + mgq \Rightarrow \frac{\partial T}{\partial t} + \frac{1}{2m} \left(\frac{\partial T}{\partial p}\right)^2 - mg \frac{\partial T}{\partial p} = 0$

(3)

c) let $T(\phi, t) = W(\phi) - Qt \Rightarrow Q = \frac{1}{2m}\phi^2 - mg \frac{dW}{d\phi}$

or $T = \frac{1}{mg} \int d\phi \left(\frac{1}{2m}\phi^2 - Q \right) - Qt \Rightarrow T = \frac{1}{6m^2g} \phi^3 - \frac{Q}{mg} \phi - Qt$

d) $P = -\frac{\partial T}{\partial Q} = \frac{\phi}{mg} + t$

$g = -\frac{\partial T}{\partial \phi} = -\frac{\phi^2}{2m^2g} + \frac{Q}{mg} \Rightarrow Q = \frac{1}{2m}\phi^2 + mgg$

Q is the energy, and P is the time when $\phi=0$ (i.e., the particle is at rest and its height is a maximum.)

$\phi = mg(P - t)$

$g = -\frac{1}{2m^2g} \phi^2 + \frac{Q}{mg} \Rightarrow g = -\frac{1}{2}g(t - P)^2 + \frac{Q}{mg}$